



Exercise 1:

Determination of the angular frequency of the system:

According to the fundamental relation of dynamics

(The equation of motion from Newton's law is)

$$\sum \vec{F} = m\vec{y}$$

The system at rest, thus in equilibrium, we have:

$$\sum \vec{F} = \vec{0} \Rightarrow \vec{P} + \vec{T} = \vec{0}$$

(tension \vec{T} or restoring force (\vec{F}_r))

$$\Rightarrow mg - k\Delta l = 0, \quad \Delta l = \text{Spring elongation}$$

$$mg = k\Delta l \Rightarrow k = \frac{mg}{\Delta l}$$

The Lagrangian : $L = T - V$

kinetic energy of the system : $T = \frac{1}{2}m\dot{x}^2$

Potential energy of the spring : $V = \frac{1}{2}kx^2$

$$\Rightarrow L = \frac{1}{2}m\dot{x}^2 - \frac{1}{2}kx^2$$

The Lagrange equation:

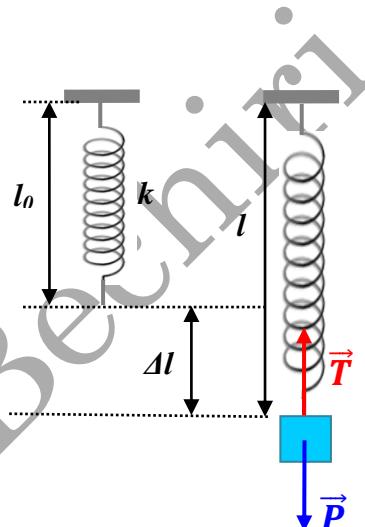
$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \left(\frac{\partial L}{\partial x} \right) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) = \frac{d}{dt} \left[\frac{\partial}{\partial \dot{x}} \left(\frac{1}{2}m\dot{x}^2 \right) \right] = \frac{d}{dt} (m\dot{x}) = m\ddot{x}$$

$$\frac{\partial L}{\partial x} = \frac{\partial}{\partial x} \left(-\frac{1}{2}kx^2 \right) = -kx$$

$$\Rightarrow m\ddot{x} + kx = 0$$

$$\Rightarrow \ddot{x} + \frac{k}{m}x = 0$$



$$\begin{cases} \ddot{x} + \frac{k}{m}x = 0 \\ \ddot{x} + \omega_0^2 x = 0 \end{cases} \Rightarrow \omega_0^2 = \frac{k}{m} \Rightarrow \omega_0 = \sqrt{\frac{k}{m}}$$

$$\text{the angular frequency } \omega_0 = \sqrt{\frac{mg}{\Delta l}} = \sqrt{\frac{g}{\Delta l}} = \sqrt{\frac{9.81}{7,787 \cdot 10^{-3}}} \\ \Rightarrow \omega_0 = 35,49 \text{ rad/s}$$

Exercise 2:

For a harmonic oscillator:

$$x(t) = A \cos(\omega_0 t \pm \varphi)$$

With A as the amplitude and the angular frequency ω_0

$$\omega_0^2 = \frac{k}{m}$$

$$E_{\text{mechanical}} = T + V = \frac{1}{2}kx^2 + \frac{1}{2}m\dot{x}^2 = \frac{1}{2}kA^2 = \frac{1}{2}m\omega_0^2 A^2$$

$$A = \sqrt{\frac{2E_m}{m\omega_0^2}} \\ A = 0,224m$$

Exercise 3:

a) The mass m of the solid.

$$\omega = 8 \text{ rad/s} \Rightarrow \omega^2 = \frac{k}{m} \Rightarrow m = \frac{k}{\omega^2} = \frac{48}{8^2} = 0.75 \text{ kg}$$

b) The amplitude of the motion..

$$E_m = E_c + E_p = \frac{1}{2}kx_m^2 = 0.24$$

$$x_m^2 = \frac{2}{k}0.24 = 0.01, x_m = 0.1m = 10cm$$

c) The maximum velocity of the oscillator.

$$x = x_m \sin(8t - \pi)$$

$$\dot{x} = 8x_m \cos(8t - \pi)$$

$$v_m = 8x_m = 80 \text{ cm/s}$$

d) The elongation of the oscillator for which the potential energy is equal to half of the kinetic energy.

$$E_m = E_c + E_p = 0.24$$

$$E_m = E_c + \frac{1}{2}E_c == 0.24$$

$$E_c + \frac{1}{2}E_c = 0.24, \quad 3E_c = 2 \times 0.24 = 0.48$$

$$E_c = \frac{1}{3}0.48 = 0.16$$

$$E_c = \frac{1}{2}mv^2 = \frac{1}{2}m(8xx_m \cos(8t - \pi))^2 = \frac{1}{2}0.75 \times 64xx_m^2(\cos(8t - \pi))^2 = 0.16$$

$$x_m = \mp 8,16 \text{ cm}$$

e) The components of velocity and acceleration at this point.

$$v_x = \dot{x} = 8x x_m \cos(8t - \pi)$$

$$v_x = \mp 8 \times 0.0816 = \mp 0.65 \text{ m/s}$$

$$\gamma_x = \ddot{x} = 64x x_m \sin(8t - \pi)$$

$$\gamma_x = \mp 64 \times 0.0816 = \mp 5.22 \text{ m/s}^2$$

Exercise 4:

$$\omega_0^2 = \frac{k}{m} = 4 \quad \text{at } t = 0, \quad x_0 = 2\text{cm}, \quad \dot{x}_0 = 8\text{cm/s}$$

The resulting motion is : $x(t) = A\cos(\omega_0 t + \varphi)$

$$\text{such that } \omega_0 = \sqrt{\frac{k}{m}} = \sqrt{4} = 2 \text{ rad/s}$$

$$\dot{x}(t) = -A\omega \sin(\omega_0 t + \varphi)$$

$$\dot{x}(0) = -A\omega_0 \sin\varphi = 8 \dots \dots \dots \quad (2)$$

$$\frac{(2)}{(1)} = \frac{-A\omega_0 \sin \varphi}{A \cos \varphi} = -\omega_0 \tan \varphi = \frac{-8}{2} = 4 \Rightarrow \tan \varphi = \frac{-4}{\omega_0} = \frac{-4}{2} = -2$$

$$\Rightarrow \varphi = -63,43^\circ = -1,10 \text{ rad}$$

$$A \cos \varphi = 2 \Rightarrow A = \frac{2}{\cos \varphi} = 4,47 \text{ cm}$$

$$x(t) = 4,47 \cos(2t - 1,10) \quad (cm)$$

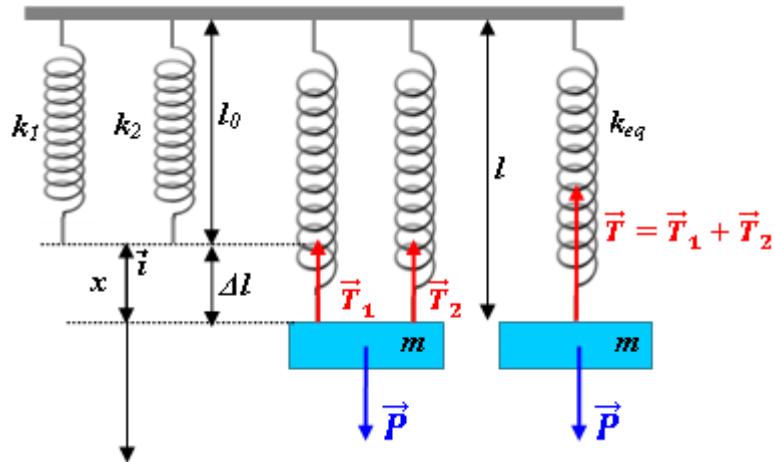
The maximum acceleration: $\ddot{x}(t) = -A\omega_0^2 \cos(\omega_0 t + \varphi)$

$$\text{Alors } \ddot{x}_{max} = |-A\omega_0^2| = 4,47 \times 2^2 = 17,88 \text{ } cms^{-2}$$

Exercise 5:

The equivalent stiffness constants k_{eq}

a) Springs in parallel :



Springs 1 and 2 are subjected to tension or restoring forces \vec{T}_1 (\vec{F}_{r1}) and \vec{T}_2 (\vec{F}_{r2}) respectively, defined by:

$$\vec{T}_1 = k_1 \Delta l \vec{i} \quad \text{et} \quad \vec{T}_2 = k_2 \Delta l \vec{i}$$

$$\Delta l = l - l_0 = x$$

At equilibrium, we have:

$$\sum \vec{F} = \vec{0} \Rightarrow \vec{P} + \vec{T}_1 + \vec{T}_2$$

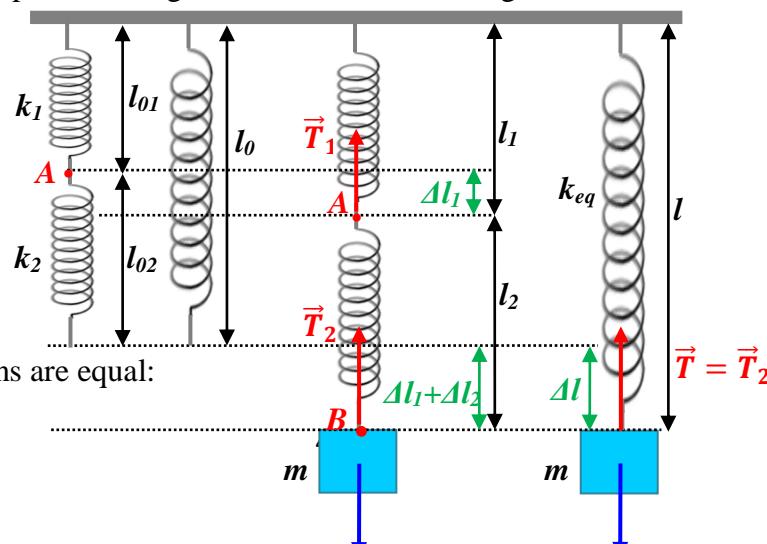
After projection onto the axis of motion, we have:

$$\begin{aligned} P - T_1 - T_2 &= 0 \\ mg - k_1 x - k_2 x &= 0 \Rightarrow mg = k_1 x + k_2 x \\ \Rightarrow mg &= (k_1 + k_2) x = k_{eq} x \end{aligned}$$

With $k_{eq} = k_1 + k_2$ for the parallel system.

b) Springs in series :

Let l_{01} , l_{02} and l_0 be the unstressed lengths of springs 1, 2, and the equivalent spring, and Δl_1 , Δl_2 and Δl their respective elongations, as shown in the figure below



At point A, the tensions are equal:

Where: $\Delta l_1 = l_1 - l_{01}$

$$\Delta l_2 = l_2 - l_{02}$$

At point B, at equilibrium, we have:

$$\begin{cases} mg = k_2 \Delta l_2 \\ \Delta l = \Delta l_1 + \Delta l_2 \end{cases}$$

According to equation (1), we find:

$$\Delta l_1 = \frac{k_2}{k_1} \Delta l_2$$

$$\Rightarrow \Delta l = \frac{k_2}{k_1} \Delta l_2 + \Delta l_2 = \left(\frac{k_2}{k_1} + 1 \right) \Delta l_2$$

$$\Rightarrow \Delta l_2 = \left(\frac{k_1}{k_1 + k_2} \right) \Delta l$$

By substituting this expression into the equilibrium equation, we obtain:

$$mg = k_2 \left(\frac{k_1}{k_1 + k_2} \right) \Delta l = \left(\frac{k_1 k_2}{k_1 + k_2} \right) \Delta l$$

$$\Rightarrow k_{eq} = \frac{k_1 k_2}{k_1 + k_2} \quad \text{for the series system}$$

$$\text{or} \quad \frac{1}{k_{eq}} = \frac{1}{k_1} + \frac{1}{k_2}$$

Analogy with equivalent capacitances C_{eq} of capacitors: it is known that for capacitors in series $\frac{1}{C_{eq}} = \sum_i \frac{1}{C_i}$ and for capacitors in parallel $C_{eq} = \sum_i C_i$

The calculation of the equivalent capacitance in an electrical oscillator is analogous to the calculation of the equivalent stiffness in a mechanical oscillator.

Exercise 6:

Determination of k_2 as a function of k_1 : ($k_1, m_1 = m$)

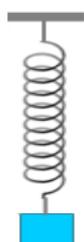
$$L = T - V$$

$$L = \frac{1}{2}m_1\dot{x}^2 - \frac{1}{2}k_1x^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \left(\frac{\partial L}{\partial x} \right) = 0$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) = \frac{d}{dt} \left[\frac{\partial}{\partial \dot{x}} \left(\frac{1}{2} m_1 \dot{x}^2 \right) \right] = \frac{d}{dt} (m_1 \dot{x}) = m_1 \ddot{x}$$

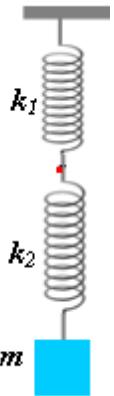
$$\frac{\partial L}{\partial x} = \frac{\partial}{\partial x} \left(-\frac{1}{2} k_1 x^2 \right) = -k_1 x$$



$$\Rightarrow m_1 \ddot{x} + k_1 x = 0$$

$$\Rightarrow \ddot{x} + \frac{k_1}{m_1} x = 0$$

$$\Rightarrow \omega_1^2 = \frac{k_1}{m_1} \Rightarrow \omega_1 = \sqrt{\frac{k_1}{m_1}} \quad \text{The natural frequency}$$



-A second spring is added in series with the first spring

$$(k_{eq}, \quad m_1) \rightarrow \omega_2 = \frac{1}{2} \omega_1$$

According to the previous answer, we find the following differential equation:

$$\ddot{x} + \frac{k_{eq}}{m_1} x = 0$$

with $\omega_2^2 = \frac{k_{eq}}{m_1} \Rightarrow \omega_2 = \sqrt{\frac{k_{eq}}{m_1}}$

The equivalent spring constant k_{eq} in the series system is:

$$k_{eq} = \frac{k_1 k_2}{k_1 + k_2}$$

$$\Rightarrow \omega_2 = \sqrt{\frac{\frac{k_1 k_2}{k_1 + k_2}}{m_1}}$$

We have : $\omega_2 = \frac{1}{2} \omega_1 \Rightarrow \omega_2^2 = \frac{1}{4} \omega_1^2$

$$\Rightarrow \frac{k_1 k_2}{k_1 + k_2} \cdot \frac{1}{m_1} = \frac{1}{4} \frac{k_1}{m_1}$$

$$\Rightarrow k_1 + k_2 = 4k_2$$

$$\Rightarrow k_1 = 3k_2$$

$$\Rightarrow k_2 = \frac{1}{3} k_3$$

Exercise 7 :

The Lagrangian of the system:

$$L = T - V$$

The kinetic energy of the system is given by:

$$T = T_M + T_m = \frac{1}{2} M v_M^2 + \frac{1}{2} m v_m^2$$

$$\vec{v}_M = \frac{d \overrightarrow{OA}}{dt} = \dot{x} \vec{i}$$

$$\begin{aligned}
&\Rightarrow T_M = \frac{1}{2} M v_M^2 = \frac{1}{2} M \dot{x}^2 \\
&\vec{v}_m = \frac{d \overrightarrow{OB}}{dt} = \frac{d}{dt} (\overrightarrow{OA} + \overrightarrow{AB}) \\
&\overrightarrow{OA} \begin{pmatrix} x \\ 0 \end{pmatrix}, \quad \overrightarrow{AB} \begin{pmatrix} l \sin \theta \\ -l \cos \theta \end{pmatrix} \\
&\Rightarrow \vec{v}_m = \frac{d}{dt} (x \vec{i} + l \sin \theta \vec{i} - l \cos \theta \vec{j}) = \\
&= \dot{x} \vec{i} + l \dot{\theta} \cos \theta \vec{i} + l \dot{\theta} \sin \theta \vec{j} = (\dot{x} + l \dot{\theta} \cos \theta) \vec{i} + l \dot{\theta} \sin \theta \vec{j} \\
&v_m^2 = x^2 + y^2 \\
&v_m^2 = (\dot{x} + l \dot{\theta} \cos \theta)^2 + (l \dot{\theta} \sin \theta)^2 = \dot{x}^2 + l^2 \dot{\theta}^2 \cos^2 \theta + 2 \dot{x} l \dot{\theta} \cos \theta + l^2 \dot{\theta}^2 \sin^2 \theta
\end{aligned}$$

We have :

$$\begin{aligned}
&\cos^2 \theta + \sin^2 \theta = 1 \\
&\Rightarrow v_m^2 = \dot{x}^2 + l^2 \dot{\theta}^2 + 2 \dot{x} l \dot{\theta} \cos \theta \\
&T_m = \frac{1}{2} m v_m^2 = \frac{1}{2} m (\dot{x}^2 + l^2 \dot{\theta}^2 + 2 \dot{x} l \dot{\theta} \cos \theta) \\
&\Rightarrow T_m = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + m \dot{x} l \dot{\theta} \cos \theta \\
&T = T_M + T_m \\
&T = \frac{1}{2} M \dot{x}^2 + \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + m \dot{x} l \dot{\theta} \cos \theta \\
&T = \frac{1}{2} (M + m) \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + m \dot{x} l \dot{\theta} \cos \theta
\end{aligned}$$

The potential energy of the system is:

we have :

$$\begin{aligned}
&V = V_M + V_m \\
&V_M = \frac{1}{2} k x^2 \\
&V_m = mgh \\
&\cos \theta = \frac{L - h}{L} \Rightarrow L - h = L \cos \theta \\
&\Rightarrow h = L - L \cos \theta = L(1 - \cos \theta) \\
&\Rightarrow V_m = mgL(1 - \cos \theta) \\
&V = V_M + V_m = \frac{1}{2} k x^2 + mgL(1 - \cos \theta) \\
&L = T - V \\
&\Rightarrow L = \frac{1}{2} (M + m) \dot{x}^2 + \frac{1}{2} m l^2 \dot{\theta}^2 + m \dot{x} l \dot{\theta} \cos \theta - \frac{1}{2} k x^2 - mgL(1 - \cos \theta)
\end{aligned}$$

Note:

This system has two independent variables x and $\theta \Rightarrow$ a system with two degrees of freedom.

To derive the differential equation, it will be necessary to apply the Lagrange equations for the two variables, that is:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \left(\frac{\partial L}{\partial x} \right) = 0 \text{ et } \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \left(\frac{\partial L}{\partial \theta} \right) = 0$$

Exercise 8:

Determination of the angular frequency of the oscillatory motion:

$$L = T - V$$

The kinetic energy :

$$T = T_{translation} + T_{rotation}$$

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} J \dot{\theta}^2$$

With J as the moment of inertia,

$$\begin{aligned} \text{We have : } x &= R\theta \Rightarrow \dot{x} = R\dot{\theta} \Rightarrow \dot{\theta} = \frac{\dot{x}}{R} \\ &\Rightarrow T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} J \frac{\dot{x}^2}{R^2} \\ &\Rightarrow T = \frac{1}{2} \left(m + \frac{J}{R^2} \right) \dot{x}^2 \end{aligned}$$

The potential energy :

$$V = \frac{1}{2} k x^2$$

$$\text{So: } L = \frac{1}{2} \left(m + \frac{J}{R^2} \right) \dot{x}^2 - \frac{1}{2} k x^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \left(\frac{\partial L}{\partial x} \right) = 0$$

$$\frac{d}{dt} \left[\frac{\partial}{\partial \dot{x}} \left(\frac{1}{2} \left(m + \frac{J}{R^2} \right) \dot{x}^2 \right) \right] = \frac{d}{dt} \left(m + \frac{J}{R^2} \right) \dot{x} = \left(m + \frac{J}{R^2} \right) \ddot{x}$$

$$\frac{\partial L}{\partial x} = \frac{\partial}{\partial x} \left(-\frac{1}{2} k x^2 \right) = -kx$$

$$\left(m + \frac{J}{R^2} \right) \ddot{x} + kx = 0$$

$$\Rightarrow \ddot{x} + \frac{k}{\left(m + \frac{J}{R^2} \right)} x = 0$$

$$\Rightarrow \omega = \sqrt{\frac{k}{\left(m + \frac{J}{R^2}\right)}}$$

For a cylinder, the moment of inertia is given by $J = \frac{1}{2}mR^2$

$$\text{So : } \omega = \sqrt{\frac{2k}{3m}} \quad \text{The natural angular frequency of the system}$$

Exercise 9:

Determination of the equation of motion:

$$L = T - V$$

The kinetic energy :

$$T = \frac{1}{2}m\dot{x}^2 + \frac{1}{2}J\dot{\theta}^2$$

$$\text{We have : } x = R\theta \Rightarrow \dot{x} = R\dot{\theta} \Rightarrow \dot{\theta} = \frac{\dot{x}}{R}$$

So

$$T = \frac{1}{2}\left(m + \frac{J}{R^2}\right)\dot{x}^2$$

-The potential energy :

$$V = \frac{1}{2}kx^2$$

-Lagrangian :

$$L = \frac{1}{2}\left(m + \frac{J}{R^2}\right)\dot{x}^2 - \frac{1}{2}kx^2$$

So, the Lagrange equations are:

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$$\frac{d}{dt}\left(\frac{\partial L}{\partial \dot{x}}\right) - \left(\frac{\partial L}{\partial x}\right) = 0$$

$$\frac{d}{dt}\left[\frac{\partial}{\partial \dot{x}}\left(\frac{1}{2}\left(m + \frac{J}{R^2}\right)\dot{x}^2\right)\right] = \frac{\partial}{\partial t}\left[\left(m + \frac{J}{R^2}\right)\dot{x}\right] = \left(m + \frac{J}{R^2}\right)\ddot{x}$$

$$\frac{\partial}{\partial x}\left(-\frac{1}{2}kx^2\right) = -kx$$

Therefore, the differential equation of motion is written as:

$$\left(m + \frac{J}{R^2}\right)\ddot{x} + kx = 0$$

$$\Rightarrow x + \frac{k}{(m + \frac{J}{R^2})} x = 0$$

With a natural angular frequency: $\omega_0 = \sqrt{\frac{k}{\frac{J}{R^2} + m}}$

The solution to the differential equation is of the form:

$$x(t) = A \cos(\omega_0 t + \varphi)$$

We have : at $t = 0s \rightarrow x(0) = 4cm$ et $\dot{x}(0) = 0$

$$x(at t = 0) = A \cos \varphi = 4cm \quad (1)$$

$$\dot{x}(t) = -A\omega_0 \sin(\omega_0 t + \varphi)$$

$$\dot{x}(0) = -A\omega_0 \sin \varphi = 0$$

$$-A\omega_0 \neq 0 \Rightarrow \sin \varphi = 0$$

$$\Rightarrow \varphi = 0$$

$$(1) \Rightarrow A \cos 0 = 4cm$$

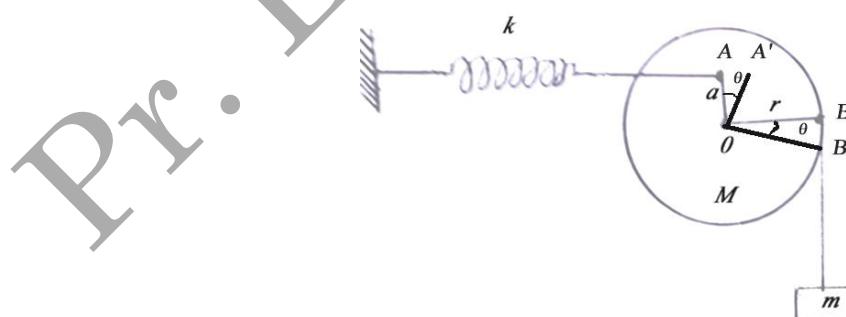
$$\Rightarrow A = 4cm$$

Therefore, we can write the equation of motion as:

$$x(t) = A \cos(\omega_0 t) = 4 \cos\left(\sqrt{\frac{k}{\frac{J}{R^2} + m}} t\right) (cm)$$

Exercise 10 :

Both the Lagrangian and energy conservation methods are equally valid and must be accepted as correct solutions



First method: Lagrangian Mechanics Analysis

Generalized Coordinate: Angular displacement θ of the disc.

The kinetic energy of the system (T):

- **Rotational kinetic energy of the disc:**

$$T_{\text{disc}} = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2$$

- **Translational kinetic energy of the mass m:**

Mass m translation at B' , $x = r\theta$, so its velocity is $\dot{x} = r\dot{\theta}$. Thus:

$$T_m = \frac{1}{2} m (r\dot{\theta})^2 = \frac{1}{2} m r^2 \dot{\theta}^2$$

- **Total kinetic energy:**

$$KE = T = T_{\text{disc}} + T_m = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} m (r\dot{\theta})^2 = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} m r^2 \dot{\theta}^2$$

Potential Energy (V):

- **Elastic potential energy of the spring:**

Spring elongation at A' , $x = a\theta$, so its potential energy is:

$$PE = V = V_s = \frac{1}{2} k (a\theta)^2 = \frac{1}{2} k a^2 \theta^2$$

- The Lagrangian is defined as:

$$L = T - V$$

Substitute T and V

$$L = T - V = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} m r^2 \dot{\theta}^2 - \frac{1}{2} k a^2 \theta^2$$

The Euler-Lagrange equation is:

- $\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = 0$

$$\frac{\partial L}{\partial \dot{\theta}}: \frac{\partial L}{\partial \dot{\theta}} = (I_{\text{disc}} + m r^2) \dot{\theta}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right): \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) = (I_{\text{disc}} + m r^2) \ddot{\theta}$$

$$\frac{\partial L}{\partial \theta}: \frac{\partial L}{\partial \theta} = -k a^2 \theta$$

Equation of Motion:

$$(I_{\text{disc}} + m r^2) \ddot{\theta} - (-k a^2 \theta) = 0$$

The equation of motion is: $(I_{\text{disc}} + m r^2) \ddot{\theta} + k a^2 \theta = 0$

Simple Harmonic Motion:

$$\ddot{\theta} + \left(\frac{ka^2}{I_{\text{disc}} + mr^2} \right) \theta = 0 \Rightarrow \ddot{\theta} + \left(\frac{ka^2}{\frac{1}{2}Mr^2 + mr^2} \right) \theta = 0 \Rightarrow \ddot{\theta} + \omega_n^2 \theta = 0$$

where: $\omega_n = \sqrt{\frac{ka^2}{I_{\text{disc}} + mr^2}}$, $\omega_n = \sqrt{\frac{ka^2}{\frac{1}{2}Mr^2 + mr^2}}$ (Natural frequency)

Second method: Energy Conservation

The kinetic energy of the system (T):

- **Rotational kinetic energy of the disc:**

$$T_{\text{disc}} = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2$$

- **Translational kinetic energy of the mass m:**

Mass m translation $x = r\theta$, so its velocity is $\dot{x} = r\dot{\theta}$. Thus:

$$T_m = \frac{1}{2} m (r\dot{\theta})^2 = \frac{1}{2} mr^2 \dot{\theta}^2$$

- **Total kinetic energy:**

$$T = T_{\text{disc}} + T_m = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} m (r\dot{\theta})^2 = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} mr^2 \dot{\theta}^2 = \frac{1}{2} (I_{\text{disc}} + mr^2) \dot{\theta}^2$$

Potential Energy (V):

- **Elastic potential energy of the spring:**

Spring elongation $x = a\theta$, so its potential energy is:

$$PE = V = V_s = \frac{1}{2} k (a\theta)^2 = \frac{1}{2} ka^2 \theta^2$$

Energy Conservation:

Total mechanical energy is constant: $T + V = \text{constant}$

Substitute T and V

$$E = T + V = \frac{1}{2} I_{\text{disc}} \dot{\theta}^2 + \frac{1}{2} mr^2 \dot{\theta}^2 + \frac{1}{2} ka^2 \theta^2 = \frac{1}{2} (I_{\text{disc}} + mr^2) \dot{\theta}^2 + \frac{1}{2} ka^2 \theta^2$$

Differentiate the energy equation with respect to time.

- $\frac{dE}{dt} = \frac{d(T+V)}{dt} = 0$
- (Applying the Chain Rule to find the derivative of a composition of functions)
- $\frac{d\dot{\theta}^2}{dt} = \frac{d\dot{\theta}^2}{d\theta} \frac{d\theta}{dt} = 2\ddot{\theta}\dot{\theta}$

$$(I_{\text{disc}} + mr^2)\ddot{\theta}\dot{\theta} + ka^2\theta\dot{\theta} = 0$$

For non-zero $\dot{\theta}$ ($\frac{d\theta}{dt} \neq 0$):

Equation of Motion:

$$(I_{\text{disc}} + mr^2)\ddot{\theta} + ka^2\theta = 0$$

Simple Harmonic Motion:

$$\ddot{\theta} + \left(\frac{ka^2}{I_{\text{disc}} + mr^2} \right) \theta = 0 \Rightarrow \ddot{\theta} + \left(\frac{ka^2}{\frac{1}{2}Mr^2 + mr^2} \right) \theta = 0 \Rightarrow \ddot{\theta} + \omega_n^2 \theta = 0$$

where: $\omega_n = \sqrt{\frac{ka^2}{I_{\text{disc}} + mr^2}}$, $\omega_n = \sqrt{\frac{ka^2}{\frac{1}{2}Mr^2 + mr^2}}$ (Natural frequency)

Exercise 11 :

$$L = T - V$$

The kinetic energy :

$$T = \frac{1}{2} m \dot{x}^2$$

-The potential energy :

$$V = \frac{1}{2} k_1 x^2 + \frac{1}{2} k_2 y^2$$

For small oscillations (small angles θ), we have

$$\theta \simeq \sin \theta = \frac{x}{b} = \frac{y}{a}$$

$$\text{So : } y = \frac{a}{b} x$$

$$V = \frac{1}{2} k_1 x^2 + \frac{1}{2} k_2 \frac{a^2}{b^2} x^2 \Rightarrow V = \frac{1}{2} \left(k_1 + \frac{a^2}{b^2} k_2 \right) x^2$$

$$L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} \left(k_1 + \frac{a^2}{b^2} k_2 \right) x^2$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \left(\frac{\partial L}{\partial x} \right) = 0$$

$$\frac{d}{dt} \left[\frac{\partial}{\partial \dot{x}} \left(\frac{1}{2} m \dot{x}^2 \right) \right] = \frac{d}{dt} (m \dot{x}) = m \ddot{x}$$

$$\frac{\partial}{\partial x} \left[- \frac{1}{2} \left(k_1 + \frac{a^2}{b^2} k_2 \right) x^2 \right] = - \left(k_1 + \frac{a^2}{b^2} k_2 \right) x$$

$$m \ddot{x} + \left(k_1 + \frac{a^2}{b^2} k_2 \right) x = 0$$

$$\Rightarrow \ddot{x} + \frac{\left(k_1 + \frac{a^2}{b^2} k_2 \right)}{m} x = 0$$

With : $\omega = \sqrt{\frac{\left(k_1 + \frac{a^2}{b^2} k_2 \right)}{m}}$ natural angular frequency